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# *Compact Phase Space*

for Loop Quantum Gravity

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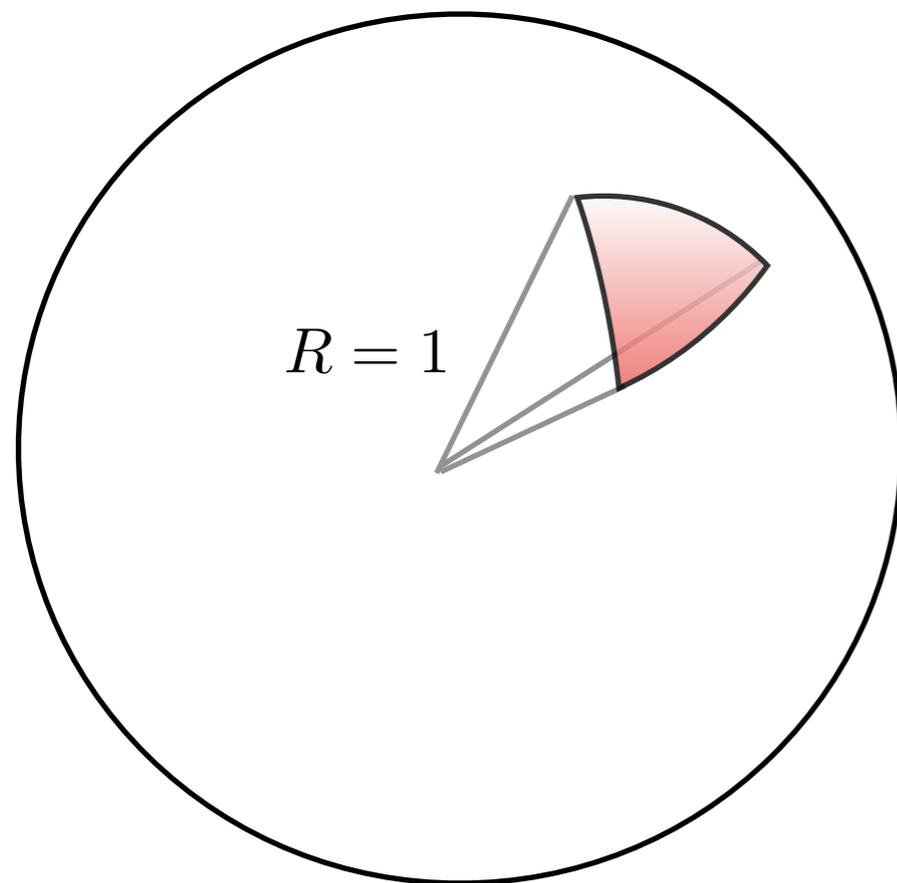
Radboud University Nijmegen



## THE PROBLEM

- Fact: our universe has a positive cosmological constant
- Does this affect the kinematics of LQG? (Dupuis-Girelli '13)
- Replacing flat cells with uniformly curved cells (Bahr-Dittrich '09)
- Classical kinematics:  $\Gamma \equiv su(2) \times SU(2) \quad \forall \ell$
- Idea: replace the algebra with the group  $\longrightarrow$  finiteness (Haggard-Han-Kaminski-Riello '14)
- Classically: compact phase space  $\longrightarrow$  finite Liouville volume
- Quantum: finite # of Planck cells, finite # orthogonal states  $\longrightarrow$  finite dim Hilbert space

$$\alpha_\ell = L_\ell / R$$



- $k_\ell$ : rotation associated to the curved arc  $\ell$
- $h_\ell$ : holonomy of the 3d connection

$$k_\ell, h_\ell \in SO(3) \sim SU(2)$$

- $SU(2) \times SU(2)$  local isometry group, or Chern-Simon gauge group

(Meusburger-Schroers '08)

- small triangles:  $k_\ell = e^{\vec{J}_\ell \cdot \vec{\tau}} \sim \mathbb{1} + J_\ell = 1 + \vec{J}_\ell \cdot \vec{\tau}$

- Standard LQG phase space:

$$SU(2) \times SU(2) \xrightarrow{R \rightarrow \infty} su(2) \times SU(2) = T^*SU(2)$$

$$(k, h) \xrightarrow{R \rightarrow \infty} (J, h)$$

- $h = e^{i\alpha}, k = e^{i\beta} \in \mathbb{C}$
- Symplectic 2-form:  $\omega = -h^{-1}dh \wedge k^{-1}dk$
- Poisson brackets:  $\{k, h\} = hk$

## ■ QUANTIZATION

- Hilbert space:  $|n\rangle \quad n = 1, \dots, N = \dim \mathcal{H}$

- Operators:  $k|n\rangle = e^{i\frac{2\pi}{N}n}|n\rangle$

$$h|n\rangle = |n+1\rangle \quad \text{cyclic: } h|N\rangle = |1\rangle$$

- Commutator:  $[h, k] = \left(e^{i\frac{2\pi}{N}} - 1\right) hk$

$$[\hat{a}, \hat{b}] = i\hbar \widehat{\{a, b\}}$$

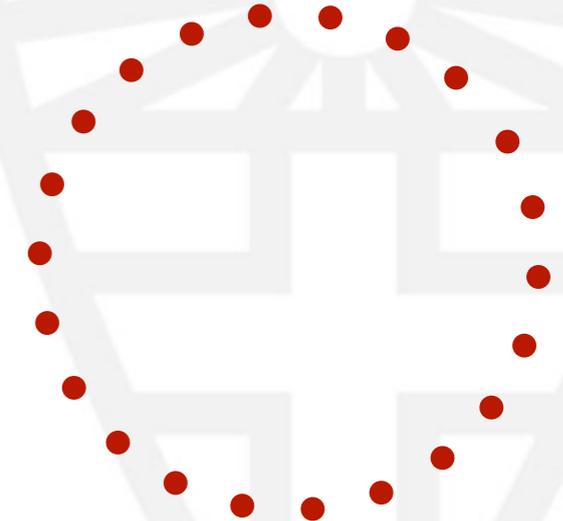
$$\hbar = \frac{2\pi}{N}$$



$$\frac{[\text{Planck length}]}{[\text{cosmological constant}]}$$

in the limit in which the radius of one of the two circles can be considered large we want to recover the symplectic form of the cotangent space

$$\omega = d\alpha \wedge d\beta$$



discrete spectrum

- $(k = e^J, h) \in SU(2) \times SU(2)$
- Symplectic 2-form:  $\omega = Tr[dk \wedge h^{-1}dh - kh^{-1}dh \wedge h^{-1}dh]$
- Poisson brackets? quasi Poisson-Lie groups

■ **QUANTIZATION**

- Hilbert space:  $L_2[SU(2)] \sim \bigoplus_{j=0}^{\infty} (\mathcal{H}_j \otimes \mathcal{H}_j)$

- Operators:  $\langle U | jmn \rangle = D_{mn}^j(U)$

- $h\psi(U) = U\psi(U)$

$$h_{AB} | jmn \rangle = \begin{pmatrix} \frac{1}{2} & j & j' \\ A & m & m' \end{pmatrix} \begin{pmatrix} \frac{1}{2} & j & j' \\ B & n & n' \end{pmatrix} | j'm'n' \rangle$$

- $J^i \psi(U) = L^i \psi(U)$

$$J^i | jmn \rangle = \tau_{mk}^{(j)} | jkn \rangle$$

limit: arc  $\ll R$   
 where  
 $\theta = Tr[kh^{-1}dh]$

- $(k = e^J, h) \in SU(2) \times SU(2)$
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limit: arc  $\ll R$   
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■ **QUANTIZATION** **QUANTUM GROUPS**

■ Hilbert space:  $L_2[SU(2)] \sim \bigoplus_{j=0}^{j_{max}} (\mathcal{H}_j \otimes \mathcal{H}_j)$

■ Operators:  $\langle U | jmn \rangle = D_{mn}^j(U)$

■  $h\psi(U) = U\psi(U)$

$$h_{AB} | jmn \rangle = \begin{pmatrix} \frac{1}{2} & j & j' \\ A & m & m' \end{pmatrix}_q \begin{pmatrix} \frac{1}{2} & j & j' \\ B & n & n' \end{pmatrix}_q | j'm'n' \rangle$$

■  $J^i \psi(U) = L^i \psi(U)$

$$J^i | jmn \rangle = \tau_{mk}^{(j)} | jkn \rangle$$

$$q^r = -1 \quad j_{max} = \frac{r-2}{2}$$

$$h_{AB}|jmn\rangle = \left( \begin{array}{ccc} \frac{1}{2} & j & j' \\ A & m & m' \end{array} \right)_q \left( \begin{array}{ccc} \frac{1}{2} & j & j' \\ B & n & n' \end{array} \right)_q |j'm'n'\rangle \text{ does not commute any more}$$

*∄ h reps*

- Wigner symbols as trivalent nodes

$$(h_{AB})_{m'n'}^{mn} = \begin{array}{c} A \\ | \\ \text{---} m \text{---} m' \\ | \\ n \text{---} n' \\ | \\ B \end{array}$$

- Acting with two operators:

$$h_{AB}h_{CD} = \begin{array}{c} A \quad C \\ | \quad | \\ \text{---} \quad \text{---} \\ | \quad | \\ B \quad D \end{array}$$

*crossing operators*

$$h_{AB}h_{CD} = R_{AC}^{A'C'} R_{BD}^{B'D'} h_{C'D'} h_{A'B'}$$

Expanding in  $\hbar$  so that  $R \sim 1 + r$  :

$$\{h_{AB}, h_{CD}\} = r_{BD}^{B'D'} h_{C'D'} h_{AB'} + r_{AC}^{A'C'} h_{C'D} h_{A'B}$$

- Inverse order:

$$h_{CD}h_{AB} = \begin{array}{c} A \quad C \\ \text{---} \quad \text{---} \\ | \quad | \\ B \quad D \end{array}$$

- We have introduced a modification of LQG kinematics
  - **compact** phase space
  - allows to introduce a **positive** cosmological constant
- **finite dimensional Hilbert space** dim determined by the ratio between the two constants:
  - quantization (physically: **Planck constant** scale)
  - simplex curvature / deformation of Poisson algebra (physically: **cosmological constant**)
- Hilbert space reduces to usual LQG one for triangles small compared to curvature radius
- A q-deformation of the dynamics:
  - renders quantum gravity finite (Turaev-Viro '92, Han '10)
  - amount to introduce the cosmological constant (Mizoguchi-Tada '91, Han '10)
- **Compactness**: discretization of the intrinsic **and extrinsic** geometry
- **Time discreteness**:  $K_{ab} \sim dq_{ab}/dt$  where  $q_{ab}(\Delta t) \sim q_{ab}(0) + dq_{ab}/dt \Delta t$   
minimum proper time Planckian, full discrete spectrum depends on cosmological constant
- Euclidean 2+1  $\rightarrow$  to do: Lorentzian 3+1 !!!

$$q = e^{i\sqrt{\Lambda\hbar G}}$$